

Ph 12b

Homework Assignment No. 8 Due: 5pm, Thursday, 11 March 2010

1. A barrier in a well (10 points).

A free quantum-mechanical particle with mass m moves inside a one-dimensional box with impenetrable walls located at $x = \pm a$. Furthermore, a *repulsive* δ -function barrier sits at the center of the well, so the potential energy function $V(x)$ in between the two impenetrable walls is given by

$$V(x) = \frac{\hbar^2}{m} \Delta \delta(x) ,$$

where $\delta(x)$ denotes the Dirac δ -function and $\Delta > 0$. As explained in Problem 2 last week, this δ -function potential causes the logarithmic derivative of the wave function $\varphi(x)$ to jump discontinuously at the origin:

$$\varphi'(0^+) - \varphi'(0^-) = 2\Delta\varphi(0) .$$

Here $\varphi'(x)$ denotes the first derivative of $\varphi(x)$, and $\varphi'(0^+)$ (respectively $\varphi'(0^-)$) denotes the limit of $\varphi'(x)$ as x approaches zero from positive (negative) values. The sign convention used here for Δ is the opposite of that used last week; now $\Delta > 0$ is the case of a repulsive barrier.

- a) For the even energy eigenstates, what is the value of the logarithmic derivative $\varphi'(x)/\varphi(x)$ at $x = 0^+$ and $x = 0^-$?
- b) For the even energy eigenstates, derive an equation that determines the wavenumber k implicitly, where $E = \hbar^2 k^2 / 2m$. Express your answer in the form

$$\Delta a = f(ka),$$

where f is a suitable function.

- c) Consider the limiting case of an infinitely strong repulsive barrier: $\Delta a \rightarrow \infty$. What are values of the energy eigenvalues in this limit, for both even and odd n ?
- d) Draw rough sketches of the wave functions for the ground state and the first excited state in the limit $\Delta a \rightarrow \infty$.

2. Reflectionless potential (15 points).

Consider a particle with mass m moving in the attractive potential

$$V(x) = -\frac{\hbar^2 k_0^2}{m} \operatorname{sech}^2(k_0 x),$$

where $\operatorname{sech}(z) = 2(e^z + e^{-z})^{-1}$ denotes the hyperbolic secant function.

- a) Show that the time-independent Schrödinger equation for this potential can be expressed as

$$\left(-\frac{d^2}{dz^2} - 2 \operatorname{sech}^2(z) \right) \varphi(z) = \bar{k}^2 \varphi(z), \quad (1)$$

where $z = k_0 x$ is a dimensionless position variable, and $\bar{k}^2 k^2 / k_0^2 = 2mE / \hbar^2 k_0^2$ is a dimensionless wavenumber.

- b) Show that

$$(i\bar{k} - \tanh(z))e^{i\bar{k}z} \quad (2)$$

solves eq.(1).

- c) Show that eq.(2) approaches $Ae^{i\bar{k}z}$ as $z \rightarrow -\infty$ and approaches $Ce^{i\bar{k}z}$ as $z \rightarrow +\infty$, where A and C are constants. What are the values of these constants.
- d) What is the transmission amplitude C/A ? Show that the transmission probability $T = |C/A|^2$ is one, and that the reflection probability $R = 1 - T$ is zero. Hence, if a wavepacket is incident on this potential from the far left, there is no reflected wave packet at all.
- e) Find an imaginary value $\bar{k} = i\bar{\kappa}$ such that $A/C = 0$ and the transmission amplitude thus diverges. For this value of $\bar{\kappa}$, there is a normalizable bound state solution, which decays exponential for both $z \rightarrow +\infty$ and $z \rightarrow -\infty$.
- f) Check that

$$\varphi(z) = \operatorname{sech}(z)$$

solves eq.(1), where $\bar{k}^2 = -\bar{\kappa}^2$ and $\bar{\kappa}$ is the value found in (e). This is the bound state solution. What is the corresponding bound state energy?

3. Bound states in a linear potential (15 points).

Consider a particle with mass m moving in the potential

$$V(x) = F|x|$$

where $|x|$ denotes the absolute value function. Thus there is a constant force F directed toward the origin.

a) Show that the time-independent Schrödinger equation for this potential can be expressed in the form (for $x \geq 0$)

$$\left(-\frac{d^2}{dy^2} + y\right) \varphi(y) = \bar{E} \varphi(y), \quad (3)$$

where

$$y = \left(\frac{\hbar^2}{2mF}\right)^{-1/3} x, \quad \bar{E} = \left(\frac{\hbar^2 F^2}{2m}\right)^{-1/3} E.$$

Equivalently, we may write

$$\frac{d^2}{dz^2} \varphi(z) = z \varphi(z)$$

where $z = y - \bar{E}$. The solution to this equation that decays as $z \rightarrow +\infty$ is the Airy function $\text{Ai}(z)$.

All real zeros of $\text{Ai}(z)$ and of its first derivative $\text{Ai}'(z)$ occur for $z < 0$. We denote the zeros of $\text{Ai}'(z)$, in order of increasing absolute value, by a_0, a_2, a_4, \dots , and we denote the zeros of $\text{Ai}(z)$ in order of increasing absolute value by a_1, a_3, a_5, \dots . These constants have the numerical values:

$$\begin{aligned} -a_0 &= 1.0188 \dots \\ -a_1 &= 2.3381 \dots \\ -a_2 &= 3.2482 \dots \\ -a_3 &= 4.0879 \dots \\ -a_4 &= 4.8201 \dots \\ -a_5 &= 5.5206 \dots \end{aligned} \quad (4)$$

- b) Show that for $n = 0, 1, 2, \dots$ there is a bound state solution to the Schrödinger equation with n nodes and dimensionless “energy” $\bar{E} = \bar{E}_n = -a_n$.

Using the WKB approximation and the connection formulas, we can derive the Bohr-Sommerfeld criterion:

$$\int_{x_1}^{x_2} dx k(x) = \pi \left(n + \frac{1}{2} \right),$$

where n is the number of nodes in the bound state wavefunction, E_n is the corresponding energy, x_1 and x_2 are the classical turning points for $E = E_n$, and $k(x)^2 = 2m(E_n - V(x))/\hbar^2$. For the harmonic potential, this WKB estimate actually agrees with the exact value of E_n , but in general there are corrections higher order in $1/n$.

- c) Apply the WKB criterion to the linear potential, deriving a formula for \bar{E}_n . For $n = 0, 1, 2, 3, 4, 5$, compare to the exact result from (b). You should find pretty good agreement for all $n \geq 1$. Furthermore, you should find (considering the odd and even values of n separately), that the agreement gets systematically better as n increases.